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## Vibrational Analysis of Electronic Transition Bands of Coronene

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The fluorescence and absorption spectra of coronene dissolved in organic solvents were measured at room temperature. Their vibrational structures were analyzed by reference to the Raman lines of crystalline coronene measured by means of a He-Ne gas laser, and the observed transition energies were compared with the calculated ones. It was suggested that the p-band of the absorption spectrum of coronene ( $23000\text{--}27000\text{ cm}^{-1}$ ) is composed of more than one electronic transition, in contrast with the case of benzene. The fluorescence and phosphorescence spectra of coronene in *n*-octane were also measured at  $77^\circ\text{K}$ , and their vibrational structures were successfully analyzed with some perturbing vibrations and a few totally-symmetric ones.

The absorption and emission spectra of the first singlet-singlet transition band of coronene in solution show distinct vibrational structures. Further, a good mirror-image correspondence between the intensities and positions of the vibrational peaks of the fluorescence spectra and those of the  $\alpha$ -band of the absorption spectra, both in solution and in vapor, was found. Therefore, it is possible to attempt the vibrational analysis of the bands; such analysis has been carried out by Sidman,<sup>1)</sup> Shpolskii,<sup>2)</sup> and Mentzel and Zimmermann.<sup>3)</sup> However, their results are somewhat tentative and insufficient because of the lack of Raman spectra and infrared absorption spectra of the hydrocarbon.

In order to analyze the vibronic structure, it is necessary to know the frequencies and symmetries of the intramolecular vibrations, which can be introduced from the observation of the Raman spectra. How-

ever, no complete measurement of the Raman spectra of coronene has yet been reported. Only one line has been found by Mentzel and Zimmermann;<sup>3)</sup> its frequency is  $1363\text{ cm}^{-1}$  in a  $\text{CS}_2$  solution and  $1346\text{ cm}^{-1}$  in a tetrahydrofuran (THF) solution. No other lines could be detected because of the poor solubility of coronene in the solvents. In this work, however, by means of a He-Ne gas laser as a light source, weak but refined Raman spectra of crystalline coronene could be obtained.

The first singlet-singlet transition band of coronene ( $23000\text{--}27000\text{ cm}^{-1}$ ), with  $\log \epsilon = 2.5$ <sup>4)</sup> in the absorption spectrum, has been interpreted as a weakly-allowed electronic transition one, as was indicated by Brocklehurst and Bowen.<sup>5)</sup> This concept is the same as for benzene, for both hydrocarbons, benzene and coronene, have a  $D_{6h}$  symmetry. For benzene, the lowest one-electron transition is orbitally fourfold-degenerate in the Hückel MO calculation. The degeneracy is removed by electron-electron interactions; then, one degenerate  $E_{1u}$  state and two nonde-

1) J. W. Sidman, *J. Chem. Phys.*, **23**, 1365 (1955).

2) E. V. Shpolskii and L. A. Klimova, *Soviet Physics, Doklady* **1**, 782 (1956).

3) N. Mentzel and H. Zimmermann, *Ber. Bunsenges. Physik, Chem.*, **67**, 582 (1963).

4)  $\epsilon$  is the extinction coefficient.

generate states,  $B_{1u}$  and  $B_{2u}$ , are given.

The absorption and fluorescence spectra of benzene vapor have been vibrationally analyzed.<sup>6)</sup> The emission spectrum of benzene arises from one quantum addition of an  $e_{2g}$  vibration of  $606\text{ cm}^{-1}$  to the symmetry-forbidden 0—0 band. The main progression is obtained by several further quanta addition of an  $a_{1g}$  vibration of  $923\text{ cm}^{-1}$ . In coronene, as was reported by Brocklehurst and Bowen,<sup>5)</sup> the vibrational structures of the fluorescence spectrum in heptane at  $-180^\circ\text{C}$  are largely composed of several series of three bands. There are three remarkable series with interpeak distances of approximately  $1360\text{ cm}^{-1}$ . This relation is common in the fluorescence spectrum and also in the  $\alpha$ -band of the absorption spectrum, so it is possible to analyze the first singlet-singlet transition bands of coronene with some  $e_{2g}$  and a few  $a_{1g}$  vibrations, as in the case of benzene.

In this paper, the first singlet-singlet transition bands of coronene will be vibrationally analyzed by reference to the Raman lines of the crystalline sample. The above transition bands and the p- and  $\beta$ -bands of the absorption spectrum will also be analyzed using the transition energies calculated by means of the molecular orbital theory.<sup>7-9)</sup> Further, the vibrational structures of the fluorescence and phosphorescence spectra of coronene in glassy *n*-octane ( $77^\circ\text{K}$ ) will be discussed.

## Experimental and Results

**1.1. Raman Spectra of Crystalline Coronene.** In order to analyze the vibronic structures of the electronic spectra of coronene, the Raman lines of crystalline coronene were measured by means of a laser Raman spectrometer with a He-Ne gas laser (Japan Electron Optics Laboratory Co.) attached. The He-Ne gas laser ( $15803\text{ cm}^{-1}$ ) was operated at an output of 40 mW. The applied sample was crystalline powder of high-purity coronene.<sup>10)</sup> A single crystal was not employed because the crystal form was of a needle-type, so the polarized Raman spectrum could not be measured.

Although a weak emission spectrum was overlapped on the Raman spectrum, several Raman lines were identified at room temperature and at the temperature of liquid nitrogen (listed in Table 1).

Because a coronene molecule has a  $D_{6h}$  symmetry, the  $e_{2g}$ ,  $e_{1g}$ , and  $a_{1g}$  vibrations should be Raman-active. Among them,  $a_{1g}$  vibrations are generally most intensive; therefore,

5) E. J. Bowen and B. Brocklehurst, *J. Chem. Soc.*, **1954**, 3875; **1955**, 4320.

6) H. Sponer, G. Nordheim, A. L. Sklar, and E. Teller, *J. Chem. Phys.*, **7**, 207 (1939).

7) N. S. Ham and K. Ruedenberg, *ibid.*, **25**, 1 (1956).

8) R. Hummel and K. Ruedenberg, *J. Phys. Chem.*, **66**, 2334 (1962).

9) H. Gutfreund and W. A. Little, *J. Chem. Phys.*, **50**, 4468 (1969); *Phys. Rev.*, **183**, 68 (1969).

10) The purification process has already been reported in another article.<sup>11)</sup>

11) S. Iwahima, K. Ohno, T. Kajiwara, and J. Aoki, *Nippon Kagaku Zasshi*, **90**, 884 (1969).

TABLE 1. RAMAN LINES OF THE CRYSTALLINE CORONENE

	RT	77°K
Lattice vibration range	40 $\text{cm}^{-1}$ 57 87	43 $\text{cm}^{-1}$ 52 77 90
Intra-molecular vibration range	S 369 $\text{cm}^{-1}$ W 435 VS 482 W 830 VS 1354 VS 1371 W 1584 W 1631	M 369 $\text{cm}^{-1}$ M 435 VS 487 M 571 W 828 S 840 VS 1354 VS 1369 W 1584 W 1625 W 1633

it can be expected that most of the strong Raman lines obtained in the intramolecular vibration range belong to  $a_{1g}$  except for one line of  $369\text{ cm}^{-1}$ ,<sup>12)</sup> since, in the lower-frequency region, in general there is no totally symmetric vibration in a large aromatic hydrocarbon such as coronene.

Among the observed lines,  $1354$  and  $1371\text{ cm}^{-1}$  are the most intensive lines in the intramolecular vibration range; they may correspond to  $1363\text{ cm}^{-1}$  measured in the  $\text{CS}_2$  solution by Mentzel and Zimmermann. The complete assignments of the Raman lines have not yet been made because of overlapping with weak emission spectra and because of the difficulty of polarized-Raman-spectra measurements.

In the lattice vibration range, there is a drastic change in the structure of the Raman lines between room temperature and the temperature of liquid nitrogen (Fig. 1). There-

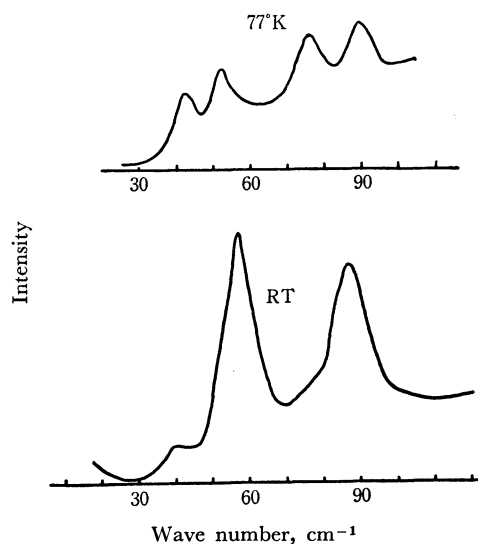


Fig. 1. The lattice vibration range of Raman spectra of coronene crystallites.

12) See Sec. 2.1.

fore, it seems that the phase transition takes place within the temperature range; the measurement of the heat capacity will now be undertaken.

**1.2. Absorption Spectra of Gaseous Coronene.** For understanding the optical properties of the coronene molecule, the absorption spectra of gaseous coronene are useful.<sup>13)</sup> This seven-membered aromatic hydrocarbon, coronene, is very stable even at high temperatures of about 500°C, so that it was possible to obtain vapour phase enough to measure the absorption spectra. Purified coronene (46 mg) was enclosed in a cylindrical cell 3 cm in diameter and 30 cm in length *in vacuo*. The cell was placed in an electric furnace and heated to a constant temperature of 400°C in order to observe the  $\alpha$ -band and to 250°C to observe the p-band. A tungsten lamp was used as the light source for the optical-absorption measurement. The most appropriate period of exposure was 20 min for the  $\alpha$ -band observation and 10 min for the p-band observation, while the slit width of the spectrograph, a Spex Model 1700 III, was 100  $\mu$ . The relative intensities of the absorption peaks of the observed  $\alpha$ -band are shown in Fig. 2.

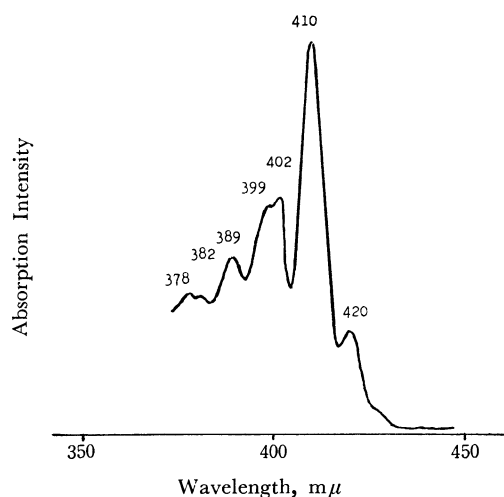


Fig. 2 The absorption spectrum of gaseous coronene.

As is shown in Table 2, the vibrational structure of the  $\alpha$ -band of gaseous coronene is very similar to that of the benzene solution of coronene at room temperature reported by Clar,<sup>14)</sup> while the p-band of the gaseous state is blue-shifted by a lack of solvent perturbation in comparison with that in solution. The fact that the  $\alpha$ -band is hardly affected at all by the solvent supports the assignment of the band to an electronic transition from the  $^1A_{1g}$  state to the  $^1B_{2u}$  state which is symmetry-forbidden. The absence of the absorption peak at about 4270 Å in the gaseous phase, in contrast with its appearance in the benzene solution, suggests that the peak corresponds to that of the symmetry-forbidden transition.

**1.3. Absorption and Emission Spectra of Coronene in Solution.** The absorption and fluorescence spectra of coronene in solution were measured at room temperature. Fig. 3 and 4 show the absorption and emission spectra of coronene in THF. The peaks are named, as is shown in the figures,

13) J. Aihara, K. Ohno, and H. Inokuchi, This Bulletin, **43**, 2435 (1970).

14) E. Clar, "Polycyclic Hydrocarbons," Academic Press Inc., N. Y. (1964), p. 83.

TABLE 2. ABSORPTION SPECTRUM OF GASEOUS CORONENE IN COMPARISON WITH THE ABSORPTION SPECTRA IN BENZENE AT ROOM TEMPERATURE

	Vapor	In benzene	In benzene by Clar <sup>14)</sup>
$\alpha$ -band	(4264 Å) <sup>a)</sup>	4275 Å	4280 Å
	4200	4205	4200
	4104	4101	4100
	4022	4028	4020
	3980	3978	3965
	3890	3883	3880
	3820	3817	3815
	3780	3782	3780
		3686	3685
p-band		3560	
		3474	
	3370	3412	3415
	3310	3359	3360
	3260	3311	
		3258	3255
		3193	3195
		3150	3165

a) Very weak absorption was observed in the neighbourhood of 4264 Å, but it is impossible to determine the position exactly.

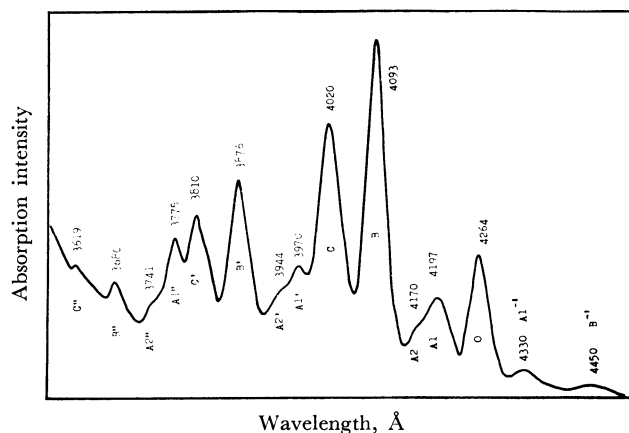


Fig. 3-1. The  $\alpha$ -band of the absorption spectrum of coronene in THF at room temperature.

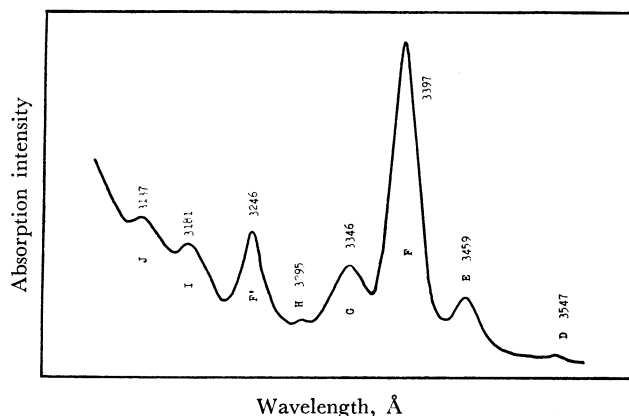


Fig. 3-2. The p-band of the absorption spectrum of coronene in THF at room temperature.

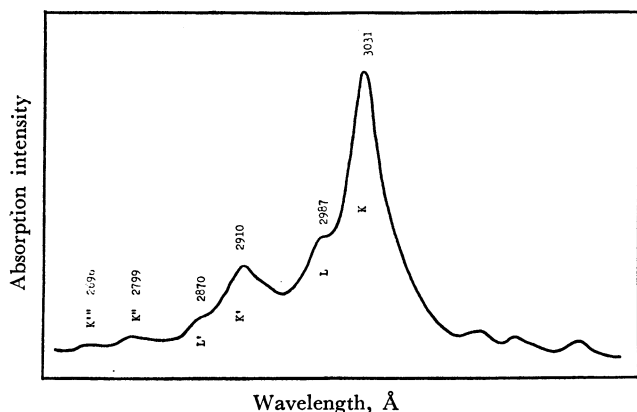


Fig. 3-3. The  $\beta$ -band of the absorption spectrum of coronene in THF at room temperature.

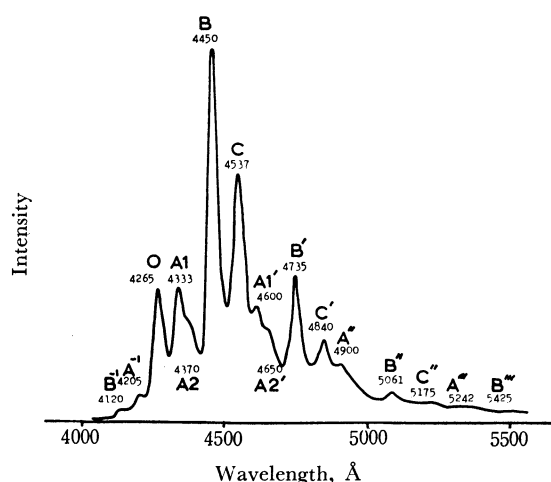


Fig. 4. The fluorescence spectrum of coronene in THF at room temperature.

according to the repeating structures and the mirror-image correspondence.

The vibrational structures of the fluorescence spectra of coronene in solution are composed of three ( $A_1$ ,  $B$ , and  $C$ ) or four series ( $A_1$ ,  $A_2$ ,  $B$ , and  $C$ ), with a common interpeak distance, as in the case of the  $\alpha$ -band of the absorption spectra of coronene. Clearly, there is a good mirror-image correspondence between the fluorescence spectra and the  $\alpha$ -band of the absorption spectra. The mirror position at about 4270 Å (O-peak) coincides with the O-O position of the first singlet-singlet transition predicted by the absorption spectrum of gaseous coronene.

**1.4. Fluorescence and Phosphorescence Spectra of Coronene in *n*-Octane at 77°K.** As is well known, the quasi-linear spectra in *n*-paraffin at low temperatures give rather fine vibronic structures.<sup>2,5</sup> Therefore, the fluorescence and the phosphorescence spectra of high-purity coronene in *n*-octane were measured at 77°K in order to analyze the vibronic structures of the  $^1A_{1g}$ - $^1B_{2u}$  transition band and also the  $^1A_{1g}$ - $^3B_{1u}$  transition one. The resolving power of the spectra observed by means of a Spex 1700 III spectrograph are satisfactory (Fig. 5); for example, the line-widths at 4320 Å and at 4440 Å are about 4 Å (21  $\text{cm}^{-1}$ ). It is noted that the envelope of this fluorescence spectrum shows a very similar pattern ( $A_1$ ,  $A_2$ ,  $B$ ,  $C$ , and their repetition) to that of the fluorescence bands in THF at room temperature. However, there exist some multiplicities. According to a recent report by Pfister,<sup>15</sup> the splittings are caused by many

possible environments of a solute molecule dissolved in *n*-paraffin crystal. Although the number of possible sites depends, in general, upon the method of crystallization, in any of our measurements there were two possible sites for the case of coronene in *n*-octane; there exist two fluorescence origins, 23529  $\text{cm}^{-1}$  and 23495  $\text{cm}^{-1}$ , but only one phosphorescence origin is observed, at 19399  $\text{cm}^{-1}$ . However, if the other phosphorescence origin is assumed to be 19411  $\text{cm}^{-1}$ , the vibronic bands can be divided into two series with similar vibronic patterns; one series starts from 19399  $\text{cm}^{-1}$ , and the other, from 19411  $\text{cm}^{-1}$ .

## Vibrational Analysis of the Spectra in THF

**2.1. First singlet-singlet Transition Bands.** The  $\alpha$ -band of the absorption spectrum shows a good mirror-image correspondence to the fluorescence spectrum. The O-peak (4264 Å in Fig. 3 or 4265 Å in Fig. 4) is in the mirror position. As has already been suggested, the O-peak corresponds to the 0—0 position of the transition between the lowest  $^1A_{1g}$  state and the lowest  $^1B_{2u}$  state, which is a symmetry-forbidden transition and which is allowed with a one-quantum addition of an  $e_{2g}$  vibration. The main progressions are obtained by the further addition of several quanta of  $a_{1g}$  vibrations.

Therefore, it is possible to analyze the vibronic structure of the first singlet-singlet transition with several  $e_{2g}$  vibrations and a few  $a_{1g}$  vibrations. The results of the analysis of the  $\alpha$ -band and the fluorescence spectrum of coronene in THF are shown in Table 3. The vibrational structure of the  $\alpha$ -band (or the fluorescence spectrum) is composed of four series ( $A_1$ ,  $A_2$ ,  $B$ , and  $C$ ), with a common interpeak distance of about 1340—1370  $\text{cm}^{-1}$ . The adding frequency  $\nu_0$  ( $a_{1g}$ ), reveals the existence of the totally

TABLE 3. VIBRATIONAL ANALYSIS OF THE  $\alpha$ -BAND AND THE FLUORESCENCE SPECTRUM OF CORONENE IN THF

	$\alpha$ -band		Fluorescence		Assignment
	$\nu$ ( $\text{cm}^{-1}$ )	$\nu - \nu_{00}$ ( $\text{cm}^{-1}$ )	$\nu$ ( $\text{cm}^{-1}$ )	$\nu_{00} - \nu$ ( $\text{cm}^{-1}$ )	
$B^{-1}$	22470	—980	24270	—820	$-\nu_3$
$A^{-1}$	23095	—355	23780	—330	$-\nu_1$
O	23450	0	23450	0	$^1A_{1g} - ^1B_{2u}$
$A_1$	23825	375	23080	370	$\nu_1(e_{2g})$
$A_2$	23980	530	22880	570	$\nu_2(e_{2g})$
$B$	24430	980	22470	980	$\nu_3(e_{2g})$
$C$	24875	1425	22040	1410	$\nu_4(e_{2g})$
$A_1'$	25190	1740	21740	1710	$\nu_1 + \nu_0(a_{1g})$
$A_2'$	25355	1905	21510	1940	$\nu_2 + \nu_0$
$B'$	25800	2350	21120	2330	$\nu_3 + \nu_0$
$C'$	26245	2795	20660	2790	$\nu_4 + \nu_0$
$A_1''$	27490	3040	20410	3040	$\nu_1 + 2\nu_0$
$A_2''$	26730	3280	20130	3320	$\nu_2 + 2\nu_0$
$B''$	27175	3725	19760	3690	$\nu_3 + 2\nu_0$
$C''$	27630	4180	19320	4130	$\nu_4 + 2\nu_0$
$A_1'''$			19080	4370	$\nu_1 + 3\nu_0$
$B'''$			18430	5020	$\nu_3 + 3\nu_0$

15) C. Pfister, *J. Chim. Phys.*, **67**, 418 (1970).

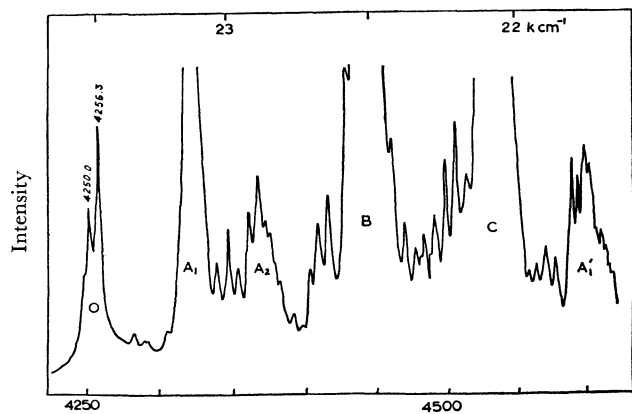


Fig. 5-1. Wavelength, Å

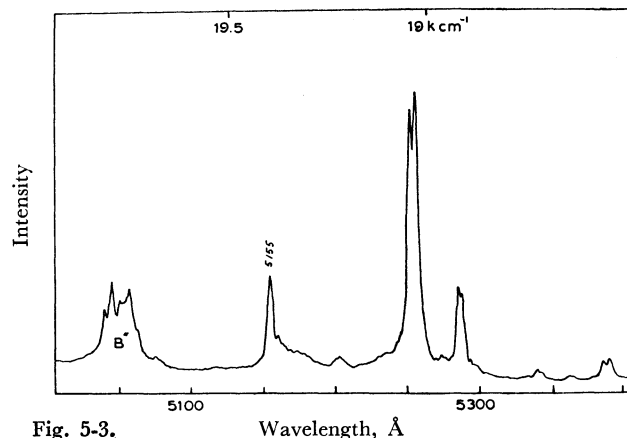


Fig. 5-3. Wavelength, Å

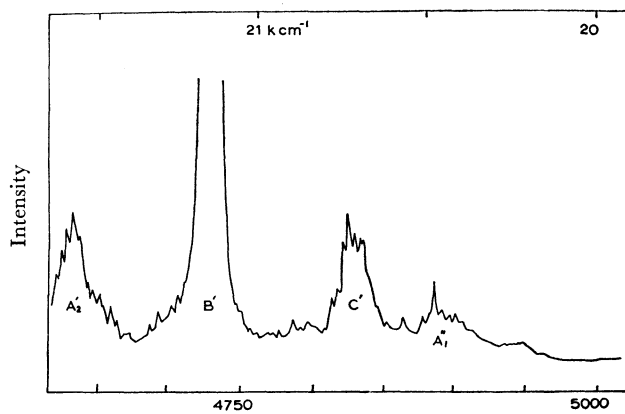


Fig. 5-2. Wavelength, Å

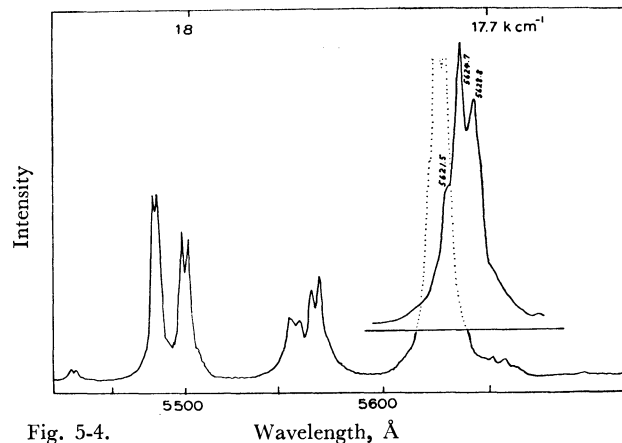


Fig. 5-4. Wavelength, Å

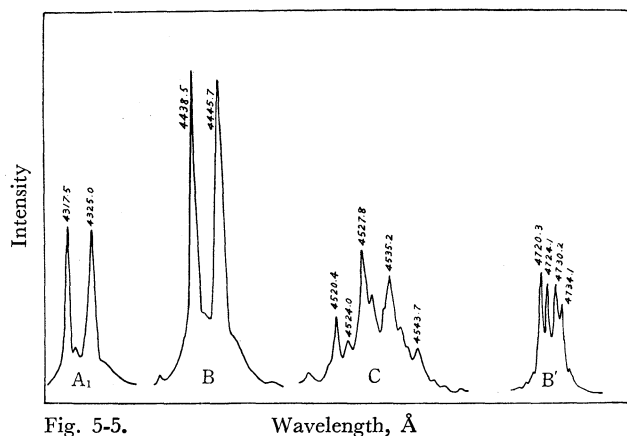


Fig. 5-5. Wavelength, Å

Fig. 5. The emission spectrum of coronene in *n*-octane at 77°K. Figure 5—5 shows the positions of peaks, A<sub>1</sub>, B, C, and B' illustrated in Fig. 5—1 and —2.

symmetric vibration of 1371 or 1354  $\text{cm}^{-1}$  found in the Raman spectrum of crystalline coronene by means of a He-Ne gas laser. This frequency also appears in both the  $\alpha$ - and  $\beta$ -bands.  $\nu_1$ ,  $\nu_2$ ,  $\nu_3$ , and  $\nu_4$  in Table 3 are perturbing vibrations belonging to  $e_{2g}$ . Among them,  $\nu_1$  (375  $\text{cm}^{-1}$ ) corresponds to the Raman line (369  $\text{cm}^{-1}$ ), which may be one of the non-totally symmetric Raman active vibrations,  $e_{2g}$  or  $e_{1g}$ , so that it is possible to predict that the Raman line (369  $\text{cm}^{-1}$ ) is an  $e_{2g}$  vibration.

Two peaks ( $A_1^{-1}$  and  $B^{-1}$  of Fig. 3—1) appearing on the longer-wavelength side of the  $\alpha$ -band, and also the two peaks,  $A_1^{-1}$  and  $B^{-1}$  of Fig. 4, on the shorter-wavelength side of the fluorescence spectrum, have

almost identical frequency differences from the O-peak which correspond well to the frequencies of the perturbing vibrations,  $\nu_1$  and  $\nu_3$ . As has already been indicated by Brocklehurst and Bowen,<sup>5)</sup> these bands seem to be the hot bands of the lowest singlet-singlet transition, for these bands could not be observed at the temperature of liquid nitrogen in either the absorption or in the fluorescence spectra. Although the agreements in the vibrational analysis between the hot-band fluorescence and the  $\alpha$ -band are not adequate, these discrepancies can easily be understood by considering the reabsorption process.

Sidman<sup>1)</sup> has assigned 990  $\text{cm}^{-1}$  (corresponding to  $\nu_3$  in our assignment) to  $\nu_1$  (370  $\text{cm}^{-1}$ ,  $e_{2g}$ ) +  $\nu$  (620

$\text{cm}^{-1}$ ,  $a_{1g}$ ) but the absence of the  $620\text{ cm}^{-1}$  line in the Raman spectrum suggests that his assignment is not correct.

**2.2.  $p$ -Band.** In the  $p$ -band of the absorption spectrum, the values of  $\log \epsilon$  range from 3 to 4.8 (Fig. 3—2). The intensity of the  $D$ -peak at  $3547\text{ Å}$  is much weaker than that of any other peak in the  $p$ -band, while it is much stronger than the peaks in the  $\alpha$ -band. Therefore, this peak seems to belong to another electronic transition.

Taking the  $E$ -peak (Fig. 3—2) at  $3459\text{ Å}$  ( $28910\text{ cm}^{-1}$ ) as the 0—0 position, we can obtain exactly the same frequencies (Table 4) as in the vibrational analysis of the  $\alpha$ -band. In the  $p$ -band, the vibration of  $530\text{ cm}^{-1}$  is the most effective perturbing vibration, and the additional frequency is  $1365\text{ cm}^{-1}$ .

The two peaks (I and J of Fig. 3—2) in the shorter-wavelength region of the  $p$ -band have an interpeak distance of  $445\text{ cm}^{-1}$ , which corresponds to the frequency of  $1420$ — $975\text{ cm}^{-1}$ , but the interpeak distance between G and I is  $1550\text{ cm}^{-1}$ , which has not yet appeared elsewhere. If I and J are the repeating structure of G and H with a one-quantum addition of a totally symmetric vibration of  $1550\text{ cm}^{-1}$ , the repetition of the most intense peak, F, must appear at the  $530 + 1550\text{ cm}^{-1}$  position. Furthermore, although H is much weaker than G, J has an intensity almost comparable to that of I. Therefore, it is unreasonable to assign I and J to the repetition of G and H. It is, then, preferable to assume the existence of another electronic transition.<sup>16)</sup> (Table 4)

TABLE 4. THE VIBRATIONAL ANALYSIS OF THE  $p$ -BAND OF CORONENE IN THF

	$\nu(\text{cm}^{-1})$	$\Delta\nu(\text{cm}^{-1})$	Assignment
D	28195		
E	28910	( 0)	$^1A_{1g} - ^1B_{1u}$
F	29440	( 530)	$\nu_5(e_{2g})$
G	29885	( 975)	$\nu_3(e_{2g})$
H	30330	(1420)	$\nu_4(e_{2g})$
F'	30805	(1895)	$\nu_5 + \nu_0(a_{1g})$
I	31435		
J	31880		

TABLE 5. THE VIBRATIONAL ANALYSIS OF THE  $\beta$ -BAND OF CORONENE IN THF

	$\nu(\text{cm}^{-1})$	$\Delta\nu(\text{cm}^{-1})$	Assignment
K	32980	0	
L	33460	480	$\nu_6$
K'	34350	1370	$\nu_0(a_{1g})$
M	34890	1910	$\nu_7$
K''	35710	2730	$2\nu_0$
K'''	37080	4100	$3\nu_0$
K''''	38460	5480	$4\nu_0$
N	39600		
P	42270		
Q	43760		
R	44480		

16) This will be considered in Sec. 3.

**2.3.  $\beta$ -Band.** The  $\beta$ -band is clearly composed of an allowed transition, with  $\log \epsilon = 5.4$  at the most intensive absorption peak, which is assigned to the  $^1A_{1g} - ^1E_{1g}$  transition, as in the case of benzene. As is shown in Table 5, the vibrational structure of the  $\beta$ -band also reveals the totally symmetric vibration of  $1370\text{ cm}^{-1}$  which appears in the  $\beta$ -band, the  $p$ -band, and the Raman spectrum.

### 3. Comparison of Transition Energies between Calculated and Observed Values

The molecular orbitals of coronene and their energies were first calculated by Moffitt and Coulson<sup>17)</sup> by means of the simple LCAO-MO method. According to their calculations, the highest occupied MO belongs to  $e_{2u}$  and the lowest vacant MO belongs to  $e_{1g}$ , so that the lowest one-electron transition is orbitally fourfold degenerate, as in the case of benzene. The electron-electron interaction removes the fourfold

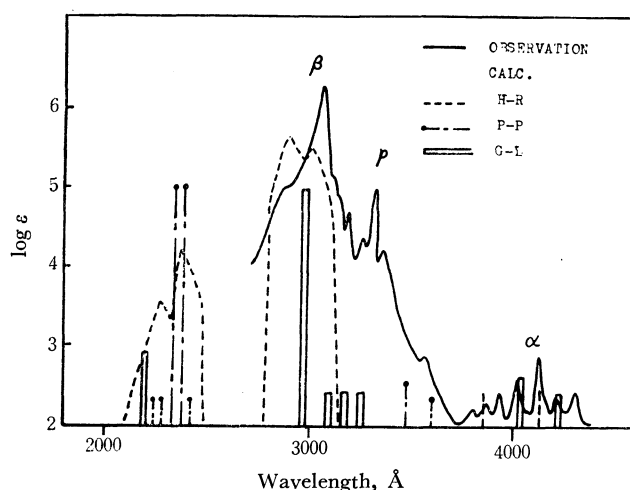


Fig. 6. The calculated and observed spectrum of coronene HR; Calculation by Hummel and Ruedenberg<sup>8)</sup> P-P; Conventional Parier-Parr MO calculation<sup>9)</sup> and G-L; Calculation by Gutfreund and Little.<sup>9)</sup> The calculated spectral lines for forbidden transitions indicate their wavelength positions only.

TABLE 6. THE TRANSITION ENERGY AND THE OSCILLATOR STRENGTH OF CORONENE<sup>a)</sup>

	$\nu(\text{cm}^{-1})$	$f$
$^1B_{2u}$	24459	0.0000
$^1B_{1u}$	27484	0.0000
$^1A_{1g}$	31751	0.0000
$^1A_{1g}$	32171	0.0000
$^1A_{1g}$	32501	0.0000
$^1E_{1u}$	33893	4.4
$^3B_{1u}$	19762	

a) In this work Nishimoto-Mataga's electron repulsion integrals<sup>18)</sup> are used as the bare Coulomb potential in the calculation of the effective interaction potential and 21 configurations are used in the CI calculation.

17) W. Moffitt and C. A. Coulson, *Proc. Phys. Soc.*, **60**, 21 (1948).

18) K. Nishimoto and M. Mataga, *Z. Physik. Chem. (Frankfurt)*, **12**, 335; **13**, 140 (1957).

TABLE 7-a. THE VIBRATIONAL ANALYSIS OF THE FLUORESCENCE SPECTRUM OF CORONENE IN *n*-OCTANE AT 77°K

	$\nu(\text{cm}^{-1})$	$\Delta\nu(\text{cm}^{-1})$	Assignment		$\nu(\text{cm}^{-1})$	$\Delta\nu(\text{cm}^{-1})$	Assignment
M	23529	0	I 0—0 ( ${}^1B_{2u}-{}^1A_{1g}$ )	W	21647	1882	II 494+1354
M	23495	34	II 0—0 ( ${}^1B_{2u}-{}^1A_{1g}$ )	W	21629	1900	II 494+1372
W	23351	178	I 178	M	21610	1919	I 570+1349
VW	23310	219	II 185	M	21591	1938	I 570+1368
VW	23295	234		W	21573	1956	II 566+1356
W	23228	301		W	21526	2003	I 632+1371
VS	23162	367	I 367	W	21503	2026	II 633+1359
W	23145	384	II 384	W	21466	2063	
VS	23121	408	II 374	W	21450	2079	
VW	23100	429	II 395	VW	21419	2110	
M	23042	487	I 487	VW	21409	2120	
M	23001	528	II 494	W	21342	2187	I 824+1363
M	22959	570	I 570	W	21313	2216	I 849+1367
M	22929	600	II 566	W	21272	2257	II 855+1368
M	22897	632	I 632	W	21252	2277	
M	22862	667	II 633	VW	21236	2293	I 947+1346
VW	22847	682	I 682	VW	21211	2318	I 947+1371
VW	22815	714	II 680	S	21184	2345	I 1001+1344
W	22766	763	I 763	S	21166	2363	I 1001+1362
VW	22730	799	II 765	S	21142	2387	II 1001+1352
M	22705	824	I 824	S	21124	2405	II 1001+1370
M	22680	849	I 849 (II 815)	VW	21102	2427	I 1071+1356
M	22640	889	II 855	VW	21073	2456	II 1071+1351
M	22582	947	I 947	VW	21053	2476	II 1071+1371
VVS	22528	1001	I 1001	W	21028	2501	I 1153+1348
VVS	22494	1035	II 1001	W	21006	2523	I 1153+1370
M	22424	1105	I 1071	W	20958	2571	I 1223+1348
M	22376	1153	I 1153	W	20940	2589	I 1223+1366
M	22336	1193	II 1159	W	20894	2635	I 1287+1348
W	22321	1208	I 1208	W	20853	2676	II 1290+1352
M	22306	1223	I 1223	W	20777	2752	I 1407+1345
W	22279	1250	II 1216	W	20760	2769	I 1425+1344
M	22274	1255	II 1221	M	20741	2788	I 1443+1345 (II 1409+1345)
M	22242	1287	I 1287	M	20725	2804	I 1457+1347
M	22205	1324	II 1290	M	20708	2821	II 1444+1343
M	22163	1366	I 1366	M	20693	2836	II 1459+1343
S	22122	1407	I 1407	M	20683	2846	I 1505+1341
S	22104	1425	I 1425	VW	20661	2868	I 1520+1348
VS	22086	1443	I 1443 (II 1409)	VW	20642	2887	II 1507+1346
S	22072	1457	I 1457 (II 1423)	W	20614	2915	II 1522+1359
VW	22056	1473		W	20565	2964	
VS	22051	1478	II 1444	W	20511	3018	
M	22036	1493	II 1459	M	20469	3060	I 367+1347+1346
M	22024	1505	I 1505	W	20448	3081	I 367+1347+1367
S	22009	1520	I 1520	W	20419	3110	II 374+1349+1353
M	21988	1541	II 1507	W	20400	3129	II 374+1349+1372
M	21973	1556	II 1522		20280	3249	
W	21956	1573	I 1573	Broad			570+2×1360
M	21930	1599	II 1565		20200	3329	
M	21901	1628	I 1628	VW	19908	3611	
M	21867	1662	II 1628	W	19841	3688	I 1001+1344+1343
M	21815	1714	I 367+1347	M	19822	3707	I 1001+1362+1344
M	21796	1733	I 367+1366	W	19798	3731	II 1001+1352+1344
M	21772	1757	II 374+1349	M	19771	3758	II 1001+1370+1353
M	21753	1776	II 374+1368	W	19701	3828	
W	21716	1813	I 1813	VW	19592	3937	
W	21701	1828	I 487+1341	VW	19535	3994	
VW	21678	1851	II 1817	W	19376	4153	I 1457+1347+1349
				VW	19361	4168	

TABLE 7-b. THE VIBRATIONAL ANALYSIS OF THE PHOSPHORESCENCE SPECTRUM OF CORONENE IN *n*-OCTANE AT 77°K

	$\nu(\text{cm}^{-1})$	$\Delta\nu(\text{cm}^{-1})$	Assignment
	(19411)	(-12)	I Assumed
S	19399	0	II ${}^3B_{1u}-{}^1A_{1g}$
M	19216	183	II 183
VS	19047	352	I 364
VS	19032	367	II 367
W	18957	442	
S	18924	475	I 487
S	18915	484	II 484
VW	18875	524	I 536
W	18758	541	II 541
W	18727	672	I 684
W	18716	683	II 683
W	18650	749	I 761
M	18570	829	I 841
M	18553	846	II 846
M	18447	952	I 964
M	18435	964	II 964
VS	18254	1145	I 1157
VS	18245	1154	II 1154
VS	18190	1209	I 1221
VS	18175	1224	II 1224
S	18005	1394	I 1406
S	17992	1407	II 1407
S	17972	1427	I 1439
S	17960	1439	II 1439
W	17789	1510	
VVS	17779	1620	I 1632
VVS	17766	1633	II 1633
W	17696	1703	I 364+1351
W	17677	1722	II 367+1355
W	17556	1843	II 484+1359
W	17289	2110	
W	16883	2516	II 1154+1362
M	16846	2553	I 1221+1344
W	16628	2771	I 1439+1344
W	16603	2796	II 1439+1357
W	16412	2987	II 1633+1354

degeneracy to give one degenerate  $E_{1u}$  state and two nondegenerate states,  $B_{1u}$  and  $B_{2u}$ .

The transition energies calculated by Ham and Ruedenberg<sup>7)</sup> and those reported by Hummel and Ruedenberg<sup>8)</sup> reproduce the experimental absorption spectra somewhat better than the conventional Pariser-Parr SCF-MO calculation, as is shown in Fig. 6.<sup>19)</sup> The agreement with the experimental results has been considered plausible. By many authors, it has been assumed that the first weak band at 23400  $\text{cm}^{-1}$  might correspond to the 2600 Å band of benzene ( ${}^1B_{2u}$ ); the band in the 29000  $\text{cm}^{-1}$  region, to the 2000 Å band of benzene ( ${}^1B_{1u}$ ), and the most intensive band, at 33000  $\text{cm}^{-1}$ , to the 1800 Å band of benzene ( ${}^1E_{1u}$ ). However, these correspondences may not

always be appropriate, because, in the case of coronene, a much larger number of electronic configurations must be considered in the CI calculation than in the case of benzene.

Recently, a more improved calculation has been presented by Gutfreund and Little.<sup>9)</sup> In their calculation, the correlation effects of the  $\pi$ -electrons were considered carefully, and it was shown that the Coulomb interaction between  $\pi$ -electrons in a conjugated molecule is screened by the presence of other  $\pi$ -electrons and that, therefore, the bare Coulomb potential must be replaced by the appropriate effective interaction potential. The obtained transition energies for the low-lying singlet states of polycyclic hydrocarbons, such as coronene, tetracene, and pyrene, agreed plausibly with the absorption spectra experimentally obtained. In the case of coronene, the ambiguous p-band can be explained clearly by their calculation. According to their calculation, it seems that the conventional assignments are not appropriate; thus, it is preferable to conclude that the  $\alpha$ -band of coronene corresponds to the two forbidden transitions; the p-band, to the three forbid-

TABLE 8. THE OBTAINED FUNDAMENTALS AND RAMAN LINES

Fluorescence		Phosphorescence		Raman
I	II	I	II	
178 $\text{cm}^{-1}$	185 $\text{cm}^{-1}$			
367	374	364	367	369
384	395			
				435
487	494	487	484	487
		536	541	
570	566			571
632	633			
682	680	684	683	
763	765	761		
824				828
849	855	841	846	840
947				
		964	964	
1001	1001			
	1071			
1153	1159	1157	1154	
1208	1216			
1223	1221	1221	1224	
1287	1290			
1347	1349			1354
1366	1368			1369
1407	(1409)	1406	1407	
1425	(1423)			
1443	1444	1439	1439	
1457	1459			
1505	1507			
1520	1522			
1573	1565			1584
1628	1628			1625
		1632	1633	1633
1813	1817			

19) The conventional Pariser-Parr MO calculation gives much larger transition energies.



den transitions, and the  $\beta$ -band, to the one allowed transition.

In the present work, we calculated the transition energies of coronene by the SCF-MO CI method, with effective electron repulsion integrals as in the Gutfreund and Little calculation. The calculated transition energies and oscillator strengths are listed in Table 6. These results show that there are three states between  ${}^1B_{1u}$  and  ${}^1E_{1u}$ , as in the results of Gutfreund and Little.<sup>9)</sup> We found that the three states (31751, 32171, and 32501  $\text{cm}^{-1}$ ) belong to  ${}^1A_{1g}$ , so that the transitions between these states and the ground states ( ${}^1A_{1g}$ ) are forbidden and vibronically allowed by  $e_{1u}$  vibrations. Although the  $e_{1u}$  vibrations are infrared-active, the infrared spectrum of coronene has not been measured; therefore, we cannot discuss further the vibronic structures of the p-band. In order to analyze the observed p-band, it would be desirable to determine the vibronic states by a calculation including the vibronic interactions among the three  $A_{1g}$  states, the lowest  ${}^1B_{1u}$  state, and the  ${}^1E_{1u}$  state. Thus, we should prefer to suggest that the p-band is composed of more than one electronic transition.

#### 4. Vibrational Analysis of Emission Spectra in *n*-Octane

The observed fluorescence and phosphorescence spectra of coronene in *n*-octane at 77°K were analyzed with reference to the results of the preliminary analysis

reported in the Experimental and Results sections and with reference to the observed Raman lines. The assignments are shown in Table 7, while the obtained fundamentals are listed in Table 8. The fundamentals obtained from the phosphorescence spectrum agree well with those from the fluorescence one; most of the obtained Raman lines are also found in the fundamentals derived from the fluorescence spectrum. The adding  $a_{1g}$  frequencies found in the fluorescence spectra are 1347  $\text{cm}^{-1}$  and 1366  $\text{cm}^{-1}$  in Series I, and 1349  $\text{cm}^{-1}$  and 1368  $\text{cm}^{-1}$  in Series II. These frequencies coincide well with the most intensive Raman lines, 1354  $\text{cm}^{-1}$  and 1369  $\text{cm}^{-1}$ . The most effective perturbing fundamentals are 367, 1001, 1407, and 1443  $\text{cm}^{-1}$  in the fluorescence spectrum and 367, 484, 1154, 1224, 1407, 1439, and 1633  $\text{cm}^{-1}$  in the phosphorescence one. As was mentioned in Sec. 2.1, almost all of these should be  $e_{2g}$  vibrations; however, 484  $\text{cm}^{-1}$ , corresponding to the intensive Raman line, 487  $\text{cm}^{-1}$ , may be an  $a_{1g}$  vibration, since the 964  $\text{cm}^{-1}$  observed in the phosphorescence spectrum approximately equals  $484 \times 2 = 968 \text{ cm}^{-1}$ .

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